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MODELING OF FLOW TRANSPORT IN POROUS MEDIA, FROM PORE SCALE TO NON-DARCY FLOW

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> Numerical modelling of porous media, Dubrovnik, September 2014



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Introduction

CO₂ sequestration & Porous media

Imaging and tomography

Flow modelling

Flow in carbonate rocks

Homogeneity

Convection term effect

Flow direction

Non-Darcy flow in Carbonates

Future works & challenges

Conclusions





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- [A] square matrix of coefficients
- d Vector between two cells centers
- K Vector in the non-orthogonality treatment
- S Source term
- S surface area vector
- t time
- V Volume
- Δ Difference
- Γ Diffusivity
- λ Under-relaxation factor
- µ Dynamic viscosity
- ρ Density
- Φ tensorial quantity

Geological Carbon Storage (GCS)

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(Adapted from original figure, courtesy of Dan Magee, Alberta Energy Utilities Board, Alberta Geologic Survey, 2008.)











Scientific Questions

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How to get more insight into fluid flow in porous media? How can we capture the complexity of porous media (specially carbonate rocks)?

Different effects in pore scale modelling?

How does system chenge when velocity increases?



Pore Scale modelling Methods



Lattice Boltzman Method (LBM)

Pore Network Model

Finite difference

Navier-Stokes equations

Finite element

Finite Volume

Advantages - Disadvantages

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Particle distribution functic

$$f_i(x + v_i, t + 1) - f_i(x, t) = \Omega_i(f_i(x, t))$$

Collision:

$$\Omega_i = -\frac{1}{\tau} (f_i - f_i^{eq}) \qquad \sum_i \Omega_i = 0 \quad and \qquad \sum_i \Omega_i v_i = 0$$

Momentum and density:

Ce

 c_3

C7

$$\sum_{i} f_i = \rho \quad and \quad \sum_{i} f_i v_i = \rho u$$

LBM

$$\int \Omega_i = 0 \quad and \qquad \sum_i \Omega_i v_i = 0$$

 c_2 C5 c_0 CI 8

C8



C4











Momentum equation

$$\rho \left(\frac{\partial v}{\partial t} + v \, . \, \nabla v \right) = - \nabla p + \mu \, \nabla^2 v$$

Continuity equation

 $\nabla . v = 0$

Expanding equations in one-direction:

$$\rho \left(\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} + w \frac{\partial u}{\partial z}\right) = -\frac{\partial p}{\partial x} + \mu \left(\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} + \frac{\partial^2 u}{\partial z^2}\right)$$
$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} = 0$$

Equation Discretisation

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Finite Volume Method



Equation Discretisation



Diffusion:

$$\int_{V} \nabla \cdot (\Gamma \nabla \phi) \, dV = \int_{S} d\mathbf{S} \cdot (\Gamma \nabla \phi) \approx \sum_{f} \Gamma_{f} \left(\mathbf{S} \cdot \nabla_{f} \phi \right)$$

For non-orthogonal meshes, correction term as below is used to preserve the second order accuracy

$$\mathbf{S} \boldsymbol{\cdot} \nabla_{\!\!f} \phi = \underbrace{|\boldsymbol{\Delta}| \nabla_{\!\!f}^{\!\!\perp} \phi}_{\text{orthogonal contribution}} + \underbrace{\mathbf{k} \boldsymbol{\cdot} (\nabla \phi)_f}_{\text{non-orthogonal correction}}$$

Where Δ and k are vectors which will be calculated by non-orthoganality treatment

Source Term:

And is linearized as

$$S_{\phi}(\phi) = \phi S_I + S_E$$

Where S_{E} and S_{f} can be dependence on $\Phi.$ The source term is integrated over control volume as

$$\int_{V} S_{\phi}(\phi) dV = S_{I} V_{P} \phi_{P} + S_{E} V_{P}$$

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SIMPLE method:

$$\frac{\partial}{\partial x_i} \left(\frac{\partial p}{\partial x_i} \right) = -\frac{\partial}{\partial x_i} \left[\frac{\partial (\rho u_i u_j)}{\partial x_j} \right]$$

Solution

Relaxation factor:

$$u^{n+1} = u^* + \alpha_u u'$$

Co. Number:

$$Co = \frac{U_f \cdot d}{\left|d\right|^2 \Delta t}$$

Boundary condition:

Fixed value:

$$S_f . (\nabla \phi)_f = \left| S_f \right| \frac{\phi_b - \phi_P}{|d|}$$

In this way, it is second order accurate if $\Phi_{\rm b}$ is constant and otherwise it's first order accurate

Fixed gradient:

$$g_b = (\frac{\mathbf{S}}{|S|} \cdot \nabla \phi)_f$$

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generic form of linear algebraic equations:

$$a_P \phi_P^n + \sum_N a_N \phi_N^n = R_P$$

Where $\Phi^n_{\ p}$ depends on the neighbouring cells

The system of algebric equations can be expressed in a matrix form of

 $[A][\phi] = [R]$

[A] ia a sparse square matrix with coefficient a_p on the diagonal and a_N off the diagonal. [Φ] is dependent variable and [R] is the source vector.

The matrix [A] can be decomposed into two matrices, the diagonal [D] and offdiagonal [N], such as:

$$[A] = [D] + [N]$$

Conjugate Gradient method has been used or solution

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Permeability
$$k = C \frac{\varphi^3}{S_v^2} = C \frac{\varphi^3}{(A_{BET}\rho_c)^2}$$
Aper is the Brunaver Emmett
Teller surface area in m²/g $A_{BET} = \frac{\sum S_n}{\rho_c N_c l_{vox}}$

K-C method

porosity of the material
$$\varphi = 1 - \frac{\langle \mu_i \rangle}{\mu_c}$$



✓ High-resolution image

- ✓ Homogeneity
- ✓ Equation effect
- ✓ Flow direction





Dierolf et al. (2010)









Change of permeability along samples

Sample Volume	Porosity	Kozeny- Carman(mD)	This work(mD)
1-1-X	17.10	20.4	24.6135
1-2-X	15.66		28.1885
2-1-X	16.62	12.6	143.6578
2-2-X	16.97		10.1517

Comparison with different methods



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Different Effects



Sample Name	Per. y (mD)	Per. z (mD)	Change y (%)	Change z (%)
1-1	33.34	23.21	35	5
1-2	10.35	24.87	63	11
2-1	3.19	17.61	97	87
2-2	11.25	8.55	10	15

Different direction effect

Sample Name	Permeability (without convective force)(mD)	Difference
1-1-X	24.5992	5.0e-4
1-2-X	28.1870	5.3e-5
2-1-X	143.6584	4.5e-6
2-2-X	10.1523	6.4e-5



Convective term effect

Pressure and Velocity field

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Non-Darcy



Extracting intrinsic permeability and non-Darcy flow parameters



Non-Darcy



Representative volume:



200³ voxel

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Darcy-Forchheimer: Darcy-Brinkmann:

Momentum:

$$-\frac{\partial \mathbf{p}}{\partial \mathbf{X}} = \frac{\mu \mathbf{v}}{\mathbf{k}} - \mu \nabla^2 \mathbf{v}$$
$$\frac{\partial}{\partial t} (\gamma \rho u_i) + u_j \frac{\partial}{\partial x_j} (\rho u_i) = -\frac{\partial p}{\partial x_i} + \mu \frac{\partial \tau_{ij}}{\partial x_j} + S_i$$

 $-\frac{\partial p}{\partial X} = \frac{\mu v}{k} + \mu \beta v^2$

Implicit method:

$$\frac{\partial}{\partial t} \int_{V} \rho \phi \, dV = \frac{\left(\rho_{p} \phi_{p} V\right)^{n} - \left(\rho_{p} \phi_{p} V\right)^{0}}{\Delta t}$$
$$\phi^{n} = \phi(t + \Delta t)$$
$$\phi^{0} = \phi(t)$$

Where Φ^n is the new value at the time step we are solving and Φ^0 is the old value

This method is first order accurate in time and guarantees the boundedness of the solution and is unconditionally stable



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Beta coefficient:

Model	equation	beta
Pascal et al.	$4.8 \times 10^{12} \times k^{-1.176}$	1343837
Coles and Hartman	$1.07 \times 10^{12} \times \varphi^{0.449} \times k^{-1.88}$	47,77
Janicek and Katz	$1.82 \times 10^8 \times \varphi^{0.75} \times k^{-1.25}$	7615,90
This work	Direct Modelling	2694,2



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REV: Representative elementary volume

Heterogeneity

Weak verification of macro results by micro results

Morphology change



Small length scale can cause high instability (high Co. number) in Navier-Stokes flow

- □ Flow direction has a significant effect on results for complicated porous media (carbonate rocks)
- □ K-C method can be applied in homogenous porous media



Acknowledgment

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Maersk Oil/Gas A/S for P³ project



CINEMA: "The Alliance for imaging of Energy Materials" the European Community for the CO2-REACT project







Thank you for your attention